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THEORY OF PROTOSTELLAR ACCRETION DISKS. S. Ruden, Department of Physics, University of California, Irvine CA 92717, USA.

I will present an overview of the current paradigm for the theory of gaseous accretion disks around young stars. Protostellar disks form from the collapse of rotating molecular cloud cores. The disks evolve via outward angular momentum transport provided by several mechanisms: gravitational instabilities, thermal convective turbulence, and magnetic stresses. I will review the conditions under which these mechanisms are efficient and consistent with the observed disk evolutionary timescales of several million years. Time permitting, I will discuss outbursts in protostellar disks (FU Orionis variables), the effect of planet formation on disk structure, and the dispersal of remnant gas.

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THERMAL CONTINUUM OF AGN ACCRETION DISKS. G. A. Shields and H. H. Coleman, Department of Astronomy, University of Texas, Austin TX 78712, USA.

We have computed the thermal continuum energy distribution of thermal radiation from the atmospheres of supermassive accretion disks around supermassive black holes, such as may power active galactic nuclei. Non-LTE radiative transfer is combined with a model of the vertical structure at each radius appropriate to the low effective gravities of these disks. Locally, the Lyman edge of H can be in emission or absorption. When the emission is summed over the disk with Doppler and gravitational redshifts taken into account, the observed continuum typically shows little sign of a discontinuity near the Lyman edge. For relatively cool disks, the Lyman edge is in absorption, but it appears as a slope change extending over several hundred angstroms, rather than an abrupt discontinuity. Disks around Kerr black holes can explain the observed range of soft X-ray luminosities of AGN, but disks around Schwarzschild holes are much too faint in soft X-rays.

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EVOLUTION OF DYNAMO-GENERATED MAGNETIC FIELDS IN ACCRETION DISKS AROUND COMPACT AND YOUNG STARS. T. F. Stepinski, Lunar and Planetary Institute, 3600 Bay Area Boulevard, Houston TX 77058, USA.

Geometrically thin, optically thick, turbulent accretion disks are believed to surround many stars. Some of them are the compact components of close binaries (X-ray binaries, cataclysmic variables), while the others are thought to be single stars (T Tauri stars). These accretion disks must be magnetized objects because the accreted matter, whether it comes from the companion star (binaries) or from a collapsing molecular cloud core (single young stars), carries an embedded magnetic field. In addition, most accretion disks are hot and turbulent, thus meeting the condition for the MHD turbulent dynamo to maintain and amplify any seed field magnetic field. In fact, for a disk's magnetic field to persist long enough in comparison with the disk viscous time it must be contemporaneously regenerated because the characteristic diffusion time of a magnetic field is typically much shorter than a disk's viscous time. This is true for most thin accretion disks. Consequently, studying magnetic fields in thin disks is usually synonymous with studying

magnetic dynamos, a fact that is not commonly recognized in the literature.

Progress in studying the structure of many accretion disks was achieved mainly because most disks can be regarded as two-dimensional flows (thin disk approximation) in which vertical and radial structures are largely decoupled. By analogy, in a thin disk, one may expect that vertical and radial structures of the magnetic field are decoupled because the magnetic field diffuses more rapidly to the vertical boundary of the disk than along the radius. Thus, an asymptotic method, called an adiabatic approximation, can be applied to accretion disk dynamo [1]. We can represent the solution to the dynamo equation in the form $B = Q(r)b(r, z)$, where $Q(r)$ describes the field distribution along the radius, while the field distribution across the disk is included in the vector function b , which parametrically depends on r and is normalized by the condition $\max |b(z)| = 1$. The field distribution across the disk is established rapidly, while the radial distribution $Q(r)$ evolves on a considerably longer timescale. It is this evolution that is the subject of this paper. The evolution of Q is dictated by the relative strength of local field amplification and radial diffusion, and is obtained numerically. Each numerical run is started from arbitrary initial conditions and is advanced in time using a numerical code based on the ISLM subroutine MOLCH.

Disks Around Compact Stars: As a first example of how a dynamo-generated magnetic field evolves in a thin accretion disk we have chosen a fiducial case of an accretion α disk around a compact star. A particular simple steady-state solution of disk structure is obtained [e.g., 2] under the assumption that the Rosseland mean opacity is approximated by Kramers' law, and radiation pressure can be neglected in comparison with gas pressure. We assume a disk surrounding a compact star of mass $M_* = 1 M_\odot$ and radius $r_* = 5 \times 10^8$ cm, with an accretion rate of 10^{16} g s⁻¹, $\alpha = 0.1$, an inner radius of $r_{in} = 2r_*$, and an outer radius of $r_{out} = 10^3 r_*$. We assume that at $t = 0$ the magnetic field is constant and has a magnitude equal to 1% of the equipartition value at the outer radius. In Fig. 1 we show the numerically calculated time evolution of the magnetic field. The nonlinearity of the dynamo equation (so-called α quenching) ensures that the magnetic field equilibrates. At first the field increases sharply at the inner radii and remains unchanged at the outer radii. By the time $t = 10^4$ s, the magnetic field in the innermost portion (up to $r = 10r_*$) of the disk achieves equilibrium. By the time $t = 10^5$ s the magnetic field in the region of the disk up to $r = 50r_*$ has reached equilibrium, and by the time $t = 10^6$ s the magnetic field in the portion of the disk within $r = 300r_*$ is in equilibrium. Finally, at $t = 10^7$ s, the magnetic field in the entire disk ($r < 10^3 r_*$) is already in equilibrium. The final magnitude of the magnetic field approaches about half of equipartition value B_{eq} . We conclude that the evolution of the magnetic field proceeds in such a way that radial transport of the magnetic field is unimportant in comparison with the local amplification, and the evolution of the magnetic field can be considered as a local phenomenon.

Disks Around Young Stars: The typical protoplanetary disk around a $1-M_\odot$ T Tauri star extends approximately from the star's surface to about 100 AU and is parameterized by $\alpha \approx 0.01$ and an accretion rate of about $10^{-6} M_\odot$ per year. At disk locations where the temperature is above about 200 K, the opacity is dominated by grains such as silicate and Fe metal grains, whereas water ice provides the dominant opacity at locations with lower temperature. In general, the temperature in the extended parts of the disk is too cool to thermally ionize the disk's gas; instead, ionization is pro-